

Contra Rb-Sr dating: an isotope fractionation mechanism for the nonradiogenic origin of excess strontium-87

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This new model for ^{87}Sr accumulation eliminates the need for either long-term decay of ^{87}Rb , or the acceleration of ^{87}Rb decay by many orders of magnitude. When diffusive isotope fractionation of Sr isotopes is repeated a sufficient number of times by within-mantle processes on isotopically-anomalous mantle xenoliths (each of which originally had low $^{88}\text{Sr}/^{86}\text{Sr}$, high $^{84}\text{Sr}/^{86}\text{Sr}$, and normal (0.7) $^{87}\text{Sr}/^{86}\text{Sr}$ ratios), the remnants of these xenoliths become isotopically heavy. In the process, these remnants also acquire near-normal (relative to the bulk Earth) $^{88}\text{Sr}/^{86}\text{Sr}$ and $^{84}\text{Sr}/^{86}\text{Sr}$ ratios, along with elevated $^{87}\text{Sr}/^{86}\text{Sr}$ ratios. When mixed with variable amounts of terrestrial nonradiogenic strontium, these remnants generate mixing lines indistinguishable from age-indicative Rb-Sr isochrons. Such mixtures also produce elevated $^{87}\text{Sr}/^{86}\text{Sr}$ signatures in oceanic crust, which are now routinely attributed, with much ambiguity, to prior crustal contamination.

Although the unreliability of the Rb-Sr method for dating is readily demonstrated,^{1,2} it is important that alternative models are developed to explain the observed patterns of strontium isotopes. The ratios of ^{88}Sr , ^{86}Sr , and ^{84}Sr , relative to each other, are commonly believed to be invariant throughout Earth materials. However, the ratio of ^{87}Sr relative to the other Sr isotopes varies widely.³ This needs understanding in the light of the creationist-diluvialist paradigm.

According to conventional uniformitarian beliefs, an elevated $^{87}\text{Sr}/^{86}\text{Sr}$ ratio (> 0.70) necessarily implies a contribution of ^{87}Sr from the decay of ^{87}Rb , the latter of which has a half-life of about 49 billion years. In fact, Rb-Sr 'isochrons' are constructed with the ratio $^{87}\text{Sr}/^{86}\text{Sr}$ as the y-axis of the

'isochron'. So how can Rb-Sr 'isochrons' be explained in the context of a young Earth? If Rb-Sr 'isochrons' are actually the result of mixing lines, the collinear relationship between $^{87}\text{Rb}/^{86}\text{Sr}$ and $^{87}\text{Sr}/^{86}\text{Sr}$ is readily explained. But what about the elevated and variable $^{87}\text{Sr}/^{86}\text{Sr}$ ratios? Some creationist scientists suppose that accelerated radioactive decay of ^{87}Rb , by many orders of magnitude, is needed to account for the existence of excess ^{87}Sr in the context of a young Earth.⁴ By contrast, in this work, a model is presented which bypasses radioactive decay of any nuclide to account for the observed $^{87}\text{Sr}/^{86}\text{Sr}$ ratios.

Isotopic fractionation

Many authors use the term 'fractionation' in a confusing manner, sometimes referring to the *partitioning* of elements, which happen to *secondarily* change isotopic ratios, or to the mixing of different sources of crustal material which *happen to have* different isotopic compositions to begin with. In proper usage, the phrase 'isotope fractionation' refers exclusively to the physical separation of one isotope of an element from another isotope of the same element.⁵

It has been known for several decades that the lighter elements (notably carbon, oxygen, hydrogen and sulphur) undergo isotopic fractionation under the influence of their physical and chemical environment, and much research has been conducted to understand this behaviour.⁶ By contrast, it has not generally been supposed that the heavier elements, such as strontium, undergo similar measurably significant fractionation on Earth.⁷

It is a sobering fact that Earth materials have not been systematically searched for isotope fractionation.⁸ In addition, isotopic fractionation is liable to be overlooked because the very process of conventional spectrometric analysis obscures it, as highlighted by Hofmann and Hart:

'Actually, the question of whether Sr isotopes can be fractionated by nature or not is irrelevant, because all Sr isotope analysis are routinely corrected for all isotopic fractionation effects, be they natural or laboratory-induced. The possibility of isotopic fractionation as an explanation for differences in $^{87}\text{Sr}/^{86}\text{Sr}$ has been proposed to both authors on many occasions. It therefore seems worthwhile to point out that isotopic fractionation does occur during the mass spectrometric ratio measurement. Strontium isotope measurements are therefore routinely corrected by measuring the (fractionated) $^{86}\text{Sr}/^{88}\text{Sr}$ ratio and comparing this value to an arbitrarily-chosen standard value of $^{86}\text{Sr}/^{88}\text{Sr}$ (=0.1194), and correcting all other measured isotope ratios (including $^{87}\text{Sr}/^{86}\text{Sr}$) using the appropriate mass dependence of fractionation. This process is inherently incapable of distinguishing between natural and mass spectrometer generated fractionation. Consequently, any variations in $^{87}\text{Sr}/^{86}\text{Sr}$ due to natural fractionation are removed by this correction.'⁹

For example, in the Allende chondrule, the effects of strontium isotope fractionation were automatically removed in the conventional normalization to $^{88}\text{Sr}/^{86}\text{Sr}$ of 8.37521,¹⁰ or its reciprocal, 0.1194, quoted above. Real, natural isotopic fractionation was later found in this meteorite only as a result of special investigation. Specialized procedures have to be followed to circumvent the ratio-normalization procedures.¹¹ The new model for heavy-isotope fractionation, developed in this work, avoids any situation where one isotopic ratio is used to correct another.

Isotopic fractionation in extraterrestrial materials

Uniformitarians believe that the elements and their respective isotopes originated from stellar nucleosynthetic processes. These nuclear processes are usually preferred over isotopic fractionation to explain isotopic anomalies, whenever isotopes of an element show a pattern that is inconsistent with mass-dependent isotope fractionation.¹²

Numerous isotope anomalies have been discovered in virtually every element studied in meteorites.^{13,14} Many of these isotopic anomalies do not fit any standard uniformitarian model of the universe's history.¹⁵

For example, the trends in calcium-isotope ratios are consistent with *mass-dependent* isotope fractionation.¹⁶ However, the trends in the calcium and magnesium anomalies are the opposite of what would result from mass dependent isotope fractionation. This prompts the rather cumbersome suggestion that a complex process of vaporization and condensation had been responsible (i.e. each element's isotopes being fractionated during separate episodes).

Furthermore, magnesium-isotope anomalies within inclusions of the Allende meteorite show a linear relationship between $^{26}\text{Mg}/^{24}\text{Mg}$ versus $^{25}\text{Mg}/^{24}\text{Mg}$,¹⁷ which is exactly what is expected from mass-dependent isotope fractionation. But there are other values for magnesium isotopes which plot far off this line, and for which the decay of now-extinct ^{26}Al is invoked.

As for strontium anomalies in meteorites, straightforward uniformitarian explanations fail once again. This is particularly true of the inferred stellar nucleosynthetic processes, most of which have difficulty accounting for the abundance of ^{88}Sr .¹⁸ Furthermore, strontium in the Murchison meteorite is isotopically light in terms of the $^{84}\text{Sr}/^{86}\text{Sr}$ ratio, yet contains a near-normal $^{88}\text{Sr}/^{86}\text{Sr}$ ratio. This is a mystery, as the same nucleosynthetic processes that generated the former ratio should also have caused the latter ratio.¹⁹

In the Allende chondrule, the $^{88}\text{Sr}/^{86}\text{Sr}$ and $^{84}\text{Sr}/^{86}\text{Sr}$ anomalies do form a simple linear regression whose negative slope is proportional to the mass differences between the respective pairs of isotopes.⁹ However, for some unknown reason, the strontium isotope anomalies are not correlated with those for magnesium. Which strontium

isotopes themselves are out of balance cannot be unambiguously answered either.²⁰

Isotopic fractionation in terrestrial materials

How common is heavy-element isotopic fractionation on Earth? It may surprise some readers to learn that there is no clear answer to this question. As noted earlier, routine studies using normalized mass spectrometric corrections obscure isotopic anomalies. Variations in the $^{86}\text{Sr}/^{88}\text{Sr}$ ratio for Earth materials commonly fall within the range of 0.1187–0.1206,²¹ in contrast to the accepted standard ratio of 0.1194.

Samples near the famous Oklo 'natural reactors' of Gabon, west Africa, have been found to have widely-altered Nd-isotope ratios,²² as well as $^{88}\text{Sr}/^{86}\text{Sr}$ ratios as high as 9.616.²³ Of course, these isotopic anomalies are caused by the local perturbation of isotopic ratios (resulting from mixture of fissionogenic strontium with bulk-Earth strontium), and not by isotopic fractionation processes. Nevertheless, one wonders whether, if isotopic anomalies are found in Sr or Nd, they will be taken as *ipso facto* evidence for past fissionogenic perturbations of isotopes.²⁴

Potassium isotopes appear to be very uniform throughout the Earth,¹³ albeit with localized exceptions.^{25,26} Natural isotopic fractionation has been recently discovered in iron,²⁷ and the same holds for terrestrial copper and zinc.²⁸

Fractionation in gases and solids

Uniformitarians believe that, with the possible exception of the earliest stages of its history, Earth materials have never been in a gaseous phase. Except for localized events such as asteroid impacts, there probably was never a circumstance during which silicate materials would ever reach their boiling point. If such reasoning is also valid within the creationist-diluvialist paradigm, then gas-phase isotope fractionation mechanisms have little relevance for the purposes of this study, and are henceforth ignored.

In like manner, diffusion-based isotopic fractionation in the solid state can be discounted. This is because the rate of diffusion in solids is so low ($\leq 10^{-16}$ cm²/sec)²⁹ that, at most, only sub-metre-sized distances could be travelled by nuclides, even over 4.5 billion years.^{30,31}

Fractionation in liquids and in electrolyte solutions

Thus, we focus our attention on the liquid phase. The element in question must diffuse from an area where it occurs at high concentration to one of low concentration (Figure 1).³² As illustrated in the figure, the isotope fractionation that occurs during diffusion is necessarily a contact phenomenon between the high- and low-concentration areas.

Since this type of isotopic fractionation was first studied in the gaseous phase, equations governing its behaviour

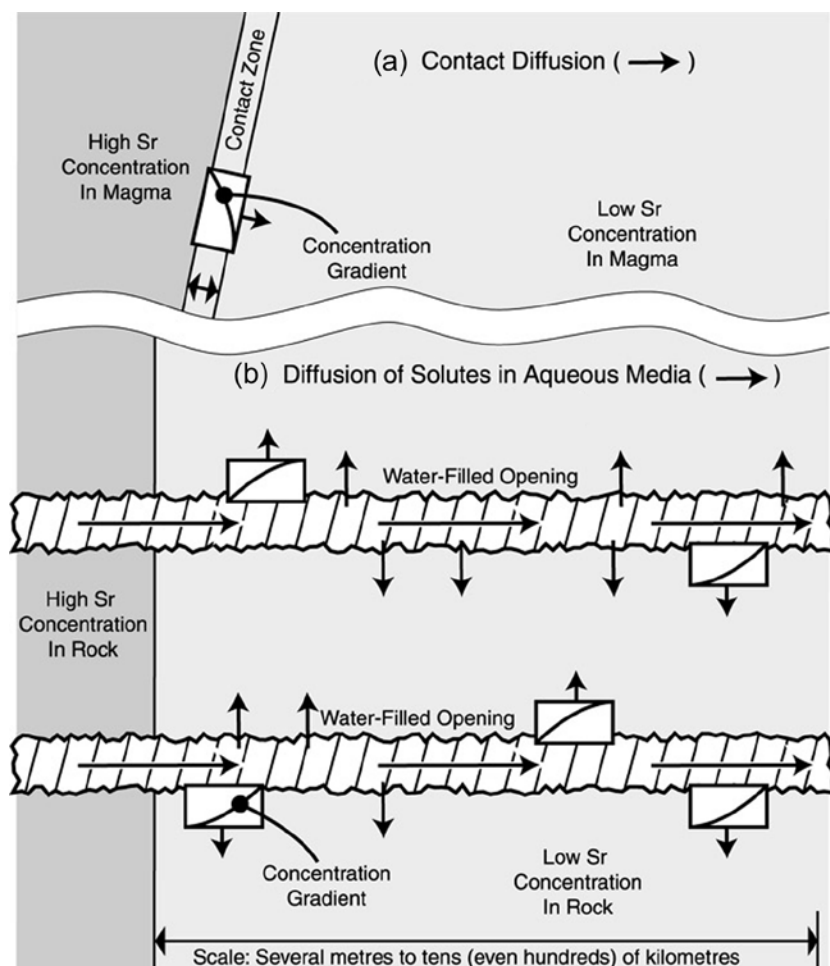


Figure 1. Diffusion of Sr in magma (a) and rock (b). The darker magma/rock to the left has a high Sr concentration while the lighter magma/rock to the right has a low Sr concentration. Diffusion of Sr occurs at the concentration gradient. In the magma (a) the concentration gradient is confined to the narrow zone at the contact, and extends only several cm at the most. For the rock (b), the concentration gradient is confined to the contact between the solution and the rock. However, the extensive porosity in the rock can allow the electrolyte solution to extend from several metres to tens (or even hundreds of kilometres).

were first described in terms of gases. The relevant equation³³ is:

$$(1) \quad D_1/D_2 = (m_1/m_2)^\beta$$

where D_1 and D_2 are the rates of diffusion, or diffusivities of isotopes of mass m_1 and m_2 , respectively, and β is a characteristic exponent, which is taken as 0.5 for a gas.³⁴

However, Richter *et al.* have performed a series of experiments which measure the degree of actual diffusion-based fractionation of calcium isotopes in simulated silicate melts (Figure 2).³⁵ From this they determined that the exponent, β , should be much lower than 0.5. Their data suggests a β value from 0.05 to 0.1, which agrees with the experiments of Tsuchiyama *et al.*³⁶ on the liquid-state diffusion of magnesium isotopes in molten MgO.³⁷

Although the diffusivity ratios of strontium isotopes in silicate melts have apparently never been measured, they should be comparable with calcium. This is because elements within the same group of the periodic table exhibit comparable diffusivities of their respective isotopes relative to each other.

The diffusion coefficients of elements in electrolyte solutions are considerably greater than those in liquid magmas. The value for the latter range from 10^{-6} to 10^{-8} cm^2/sec .^{26,30,32,33} Those in aqueous electrolyte solutions can be at least two orders of magnitude greater.³⁸ For instance, at 25°C , they are about 10^{-5} cm^2/sec ,³⁹ and this value rises to about 10^{-4} at temperatures of 500°C to 700°C .

Transport of ions in aqueous solutions greatly facilitates the potential extent of diffusive isotope fractionation in three important ways. First, the effective contact area (where an element can diffuse from a place of high concentration into one of low concentration) is greatly increased. As shown in Figure 1 (compare 1a and 1b), this follows from the greater perpendicular distance into which water can carry the material as compared with a simple non-crenellated contact. Second, the *amount* of material which can be moved in solution (Figure 1b) is much greater than can be involved in travel across a narrow interface (Figure 1a).⁴⁰ Third, the high rates of diffusion in aqueous solutions allow more material to diffuse in less time, even if the other two factors did not exist.⁴¹

All of the foregoing factors can be summarised in the standard diffusion equation:²⁷

$$(2) \quad X = (D * T)^{0.5}$$

where X is the distance diffused by a molecule, ion or isotope, D is the diffusivity, and T is the time elapsed.

Geologic constraints on isotope fractionation

Thus far, the only studies on isotopic fractionation in medium to heavy elements, in a well-defined geologic context, have been limited to contact effects in magmas, as is depicted in Figure 1a. Because of fractionation during diffusion, the element is slightly isotopically heavier on the high-concentration side of the boundary, and slightly isotopically lighter on the low-concentration side of the

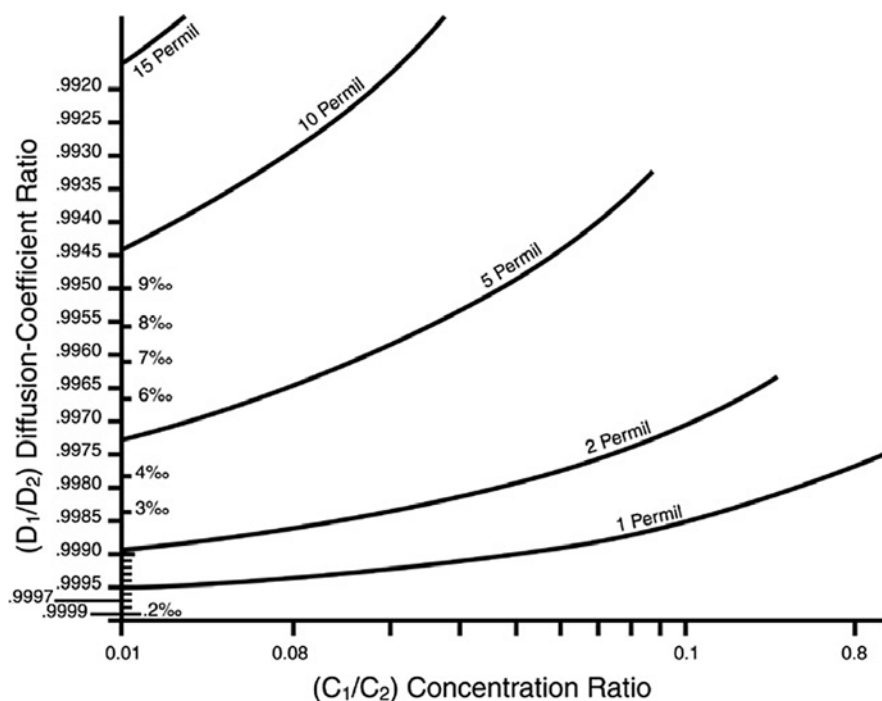


Figure 2. Degree of one-time isotopic fractionation in silicate melts (modified from Richer et al.)⁵⁸

boundary. As shown in Figure 1a, a concentration gradient exists across the boundary. With time, the gradients, shown as a hyperbolic curves throughout Figure 1, flatten out.³⁰ In other words, as noted earlier,³² the gradient is smoothed over. The previously-existing sharp difference in elemental concentrations across the magma boundaries are erased, and further isotopic fractionation of that element is thereby prevented from taking place across that boundary.

Let us now consider some examples of magma-contact isotope fractionation as observed in the solidified rocks. A lithium isotope fractionation (${}^7\text{Li}/{}^6\text{Li}$) of -30 permil was observed 6 mm into a gneiss-amphibolite contact zone, along with a -25 permil fractionation of potassium isotopes (${}^{39}\text{K}/{}^{41}\text{K}$) existing just 2 mm within the same juncture.²² Elsewhere, a granite-amphibolite contact zone of 3 cm showed a fractionation of potassium isotopes amounting to -30 permil.²³ Fractionation in other geologic environments is theoretically possible.⁴²

Regardless of the exact mechanism involved, the magnitude of isotopic fractionation, in relation to the magnitude of the concentration of the isotopes, is almost always small.⁴³ However, this conventional reasoning presupposes that isotopic fractionation in heavy elements is, at most, a one-time event. But, like the proverbial final straw that broke the camel's back, a trivial occurrence becomes significant if duplicated a sufficient number of times. In like manner, were isotopic fractionation repeated enough times, a major alteration in the isotopic ratios of the chemical elements would result. Moreover, were the repetitive isotopic fractionation on the same material to occur on a *volumetric scale* commensurate with a large fraction of the Earth's

mantle and/or crust, then appreciable amounts of the chemical elements with considerably-altered isotopic ratios would result.

With the notable exception of the repeated fractionation of hydrogen and oxygen isotopes which occur during repeated cycles of precipitation and evaporation,⁵ it does not appear that repeated fractionation of isotopes, let alone that which involves large fractions of Earth materials, has ever been contemplated, let alone researched. It is for this reason that this paper has been written. The model presented relies on the repeated movement of within-earth materials past each other (Figure 3), constantly creating new interfaces (such as shown one-time in Figure 1) through which isotopes can diffusive repeatedly, and undergoing repeated isotopic fractionation.

Large-scale fractionation of strontium isotopes

Even though God created the Universe, including the Earth, by Divine fiat on the first day (Genesis 1), we must consider the possibility of extensive movement of within-Earth materials on the first day or so of its existence as a result of speeded-up quasi-natural geochemical processes.⁴⁴ The model which I propose allows for God having assembled our planet from previously-made chunks or 'blobs' of material having divergent strontium-isotope ratios. Some of these chunks were embedded deep within the Earth, and gradually experienced equilibration of their strontium isotopes with those of the rest of the Earth through the process of isotope fractionation. In doing so, variable and often elevated ${}^{87}\text{Sr}$ abundances, relative to the other strontium isotopes, were generated. More on this shortly.

Any appreciable fractionation of strontium isotopes must be caused by a large-scale mechanism that is external to the immediate fractionation-front itself. We must consider

1. a means for relatively rapid transport of chemical species within the mantle,
2. potential mechanisms for creating and maintaining a gradient between strontium-rich and strontium-poor rocks, and
3. a means of constantly renewing this gradient.

Let us focus on aqueous transport (Figure 1b) within the lower crust and especially the mantle. The conventional belief is that, except for the first several kilometres of depth below the surface, the interior of our planet is virtually impermeable to fluids. However, there is a steadily accumulating body of evidence to the contrary.⁴⁵

Recent work by Dreibus, Jagoutz and Wanke indicates that the mantle contains about 400 ppm of water,⁴⁶ which is considerably more than generally supposed. Further evidences of water in the mantle are discussed and documented elsewhere.⁴⁷ Of course, the foregoing need not be an all-or-none situation. Parts of the Earth's surface are very wet, and others are completely dry. The same can hold for the mantle. Because of the likelihood of appreciable water in the lower crust and mantle, the capability of aqueous fluids to carry dissolved strontium to depths of several tens of kilometres (at least), and laterally therein for at least comparable distances, should be seriously considered.

Next, we must consider mechanisms, which constantly renew the concentration gradients of strontium. If there were any kind of physical movement within the mantle, as illustrated in Figure 3, it would set the stage for repeated isotope fractionation. Lamella 2 moves up (relative to Lamella 1), and the strontium diffuses across the interface near 1Z/2R. Lamella 3 moves up, in analogous fashion, relative to Lamella 2. This allows the isotopically-light Sr to diffusive across the 2R/3K boundary. Successive movements of other Lamellae allow for continuous diffusion across the boundaries marked 3K/4F, and finally, 4F/5A. The strontium is diffusing left to right in a straight line. The relative movement of Lamellae creates the *illusion* of a stepwise movement. The end result is a very isotopically-light strontium in Lamella 5.

Meanwhile, the Sr source in Lamella 1, pictured at the far upper left of Figure 3, acquires a new interface with Lamella 2 every time that Lamella 2 moves up relative to Lamella 1. With each such relocation, the Sr source preferentially sheds its light isotopes into each new contact zone material within Lamella 2. The end result is a Sr source that is very isotopically heavy.

But what kind of lamellar interfaces, as pictured in Figure 3, could actually exist in the mantle? One possibility is the sharp boundary between magmas of very different compositions. However, turbulent transport would tend to erase any incipient isotope fractionation in such circumstances. The lamellae could instead be tectonic in nature. But this seems unlikely, since the mantle material is not brittle but—to the contrary—mechanically ductile. Yet it is this very ductility which provides the answer. Patches or 'blobs' of high-Sr material could be smeared by both normal and shear strains within the mantle. This process is similar to what is believed to take place in accordance with the marble-cake model of the mantle,²⁷ except that it is a Sr-rich basaltic-composition material that gets smeared out instead of pyroxenite.

For the foregoing repeated-dilution mechanism to work (Figure 3), the surrounding mantle must have a very

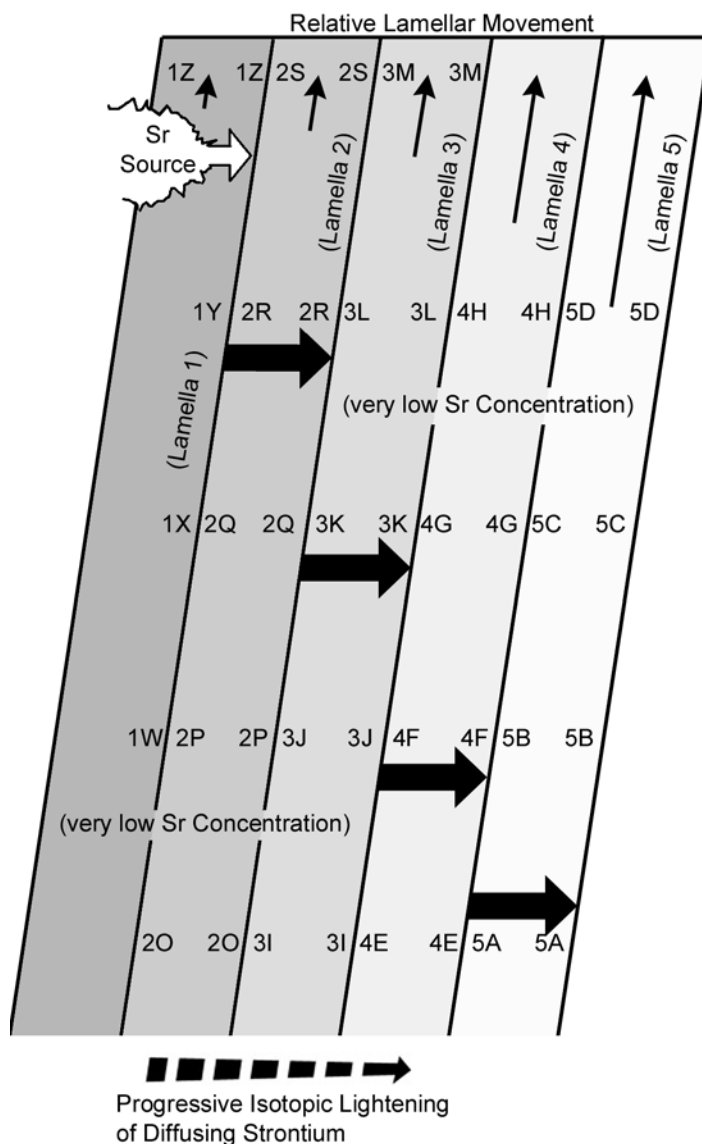


Figure 3. A mechanism which could have constantly renewed the concentration gradients of strontium within the mantle. Sr diffuses from the Sr source at the far upper left across the low-Sr lamellas, fractionating the Sr isotopes and gradually reducing the concentration gradient. Subsequently, vertical movement of the lamellas again brings low-Sr material into contact with the high Sr source allowing for multiple fractionation. The scale is comparable with that for Figure 1(b).

low (almost zero) background concentration of strontium. Such is the case. Estimates for the strontium content of the Bulk Silicate Earth (crust plus mantle) range from 10 ppm to nearly 28 ppm.⁴⁸ (Of course, the present mantle, having undergone severe extraction of strontium during the formation of the crust, has much less strontium.) This means that, were the 'blobs' of highly-concentrated strontium embedded within the pre-differentiated Earth, each of these 'blobs' would necessarily have been surrounded by a large volume of strontium-poor mantle material into which the strontium could be transported by diffusion in aqueous solutions. The repeated dilutions of strontium

contents, caused by the movements of lamellae past each other (Figure 3), would not erase the concentration gradients for a long time.

Isotopic dynamics of repeatedly-fractionated strontium isotopes

To what extent can strontium isotopes fractionate? It is well known that strontium behaves similarly to calcium, and it is recognised that calcium may at least theoretically be fractionated by interstitial fluids.⁴⁹ Let us consider the consequences of fractionating the presently-existing strontium isotopes. Using β exponent values of 0.1 and 0.05 in Equation (1), a spreadsheet program was used to calculate the isotopic abundances, which would result from the repeated diffusive isotopic fractionation of strontium isotopes. The results are presented in Tables 1 and 2.

The mass ratios of the starting points in Tables 1 and 2 are based on exact atomic mass values for each strontium isotope, taken from Holden.⁵⁰ Using Figure 1, a concentration gradient of at least 0.01 was assumed to exist for each step in the diffusive isotope fractionation process. The calculated fractionation factors, based on the atomic-mass differences of about 4, 2, and 1 between the indicated strontium isotopes, and a β value of 0.1, are 1.0084, 1.0045, and 1.0020, respectively. The corresponding fractionation factors, applicable to a β value of 0.05, are 1.0045, 1.0020, and 1.0012, respectively. Bulk-Earth Sr isotopic ratios⁴⁵ were at first used as the starting point for the fractionation experiments, even though the lowest measured $^{87}\text{Sr}/^{86}\text{Sr}$ ratio on Earth is 0.700510.⁵¹

With repeated fractionation, one gets almost pure ^{88}Sr , or, by fractionating bulk strontium in the opposite direction, almost pure ^{84}Sr (Table 1). Furthermore, as the $^{87}\text{Sr}/^{86}\text{Sr}$ ratio becomes increased above 0.7 as a result of isotope fractionation in Earth materials (Table 1), the abundance of ^{88}Sr quickly rises to nearly 100%, and that of ^{84}Sr even more quickly approaches zero. Because this type of isotopic fractionation is entirely mass-dependent, it is impossible to elevate the $^{87}\text{Sr}/^{86}\text{Sr}$ ratios without simultaneously changing all of the other strontium isotope ratios.

What about mixtures of differently-fractionated strontium? A series of mixing simulations demonstrates that isotopic fractionation is insensitive to mixing. For example, as can be seen in Table 2, the isotopic abundance after 100 continuous episodes of fractionation are almost identical to those resulting from a mixture of equal volumes of materials fractionated from a range of from zero to 200 times. This is a significant result, as it simplifies the overall analysis of complex isotopic fractionation in the following manner: One number (the times an element was fractionated) can serve as a stand-in for a series of fractionations.

However, as can be seen in Table 3, mixtures of previously-fractionated strontium cannot be combined in any manner that elevates the $^{87}\text{Sr}/^{86}\text{Sr}$ ratios and simultaneously recovers the near-normal bulk-Earth $^{86}\text{Sr}/^{88}\text{Sr}$ and $^{84}\text{Sr}/^{88}\text{Sr}$

ratios. To investigate this further, more complex hypothetical mixing simulations were performed. This time, there were repeated rounds of isotopic fractionation of strontium, in both the light and heavy directions, successively imposed upon *mixtures* of normal, isotopically-light, and/or isotopically-heavy strontium. In each case, the goal was the normal bulk-Earth $^{86}\text{Sr}/^{88}\text{Sr}$ ratio of 0.1194. A few of the results are shown in Table 4. It turns out that, under such conditions, the elevated $^{87}\text{Sr}/^{86}\text{Sr}$ ratio not only disappears, but becomes smaller than the bulk-Earth $^{87}\text{Sr}/^{86}\text{Sr}$ ratio of 0.7. One can conclude that the $^{87}\text{Sr}/^{86}\text{Sr}$ ratio can be lowered, but not elevated, by a complex alternating series of mixings and fractionations.

Up to now, all of the hypothetical fractionation simulations had some variant of bulk-Earth strontium isotopic ratios as the starting point. To explore other possibilities, a study was made of the strontium isotope ratios that are believed possible as a result of nuclear processes. While some of these ratios are comparable to those of bulk-Earth materials, others are very different.⁵² For instance, some modelled nucleosynthetic processes allow for high abundances of ^{84}Sr at the expense of all the other strontium isotopes,⁵³ while other nucleosynthetic models emphasize that ^{88}Sr can form at rates considerably less than is commensurate with its high relative abundance in bulk-Earth materials.⁵⁴ Both tendencies are shown, to varying degrees, in the isotopic starting points of the three fractionation simulations in Table 5. As for the relative abundances of ^{86}Sr and ^{87}Sr , virtually every nucleosynthetic model surveyed suggests that the former is almost always at least slightly more abundant than the latter.⁵⁵ This apparent fact is reflected in the choice of ratios of the two isotopes, relative to each other, as the starting point of each of the three hypothetical fractionating simulations shown in Table 5.

Origin of Earth materials of differing strontium isotope composition

Earlier, the consequences of isotopically-inhomogeneous masses of material within the Earth had been discussed. Also, as noted earlier, I postulate that God had created the Earth out of previously-made chunks of materials, many of which *differed* from each other, in terms of strontium-isotope concentration. (I henceforth refer to the 'blobs' as extraterrestrial xenoliths, not because they came from space, but because their isotopic composition is foreign relative to the bulk of Earth material, which had originated from other God-directed nucleosynthetic processes). I suggest that these xenoliths have been incorporated within the Earth's mantle during the time that the Earth was being assembled by God on the first day of Creation Week.

In accordance with some of the nucleosynthetic processes discussed earlier, consider a situation wherein the extraterrestrial xenoliths started out with simultaneous lower abundances of ^{88}Sr and higher abundances of ^{84}Sr than is true of bulk-Earth materials. This means that the

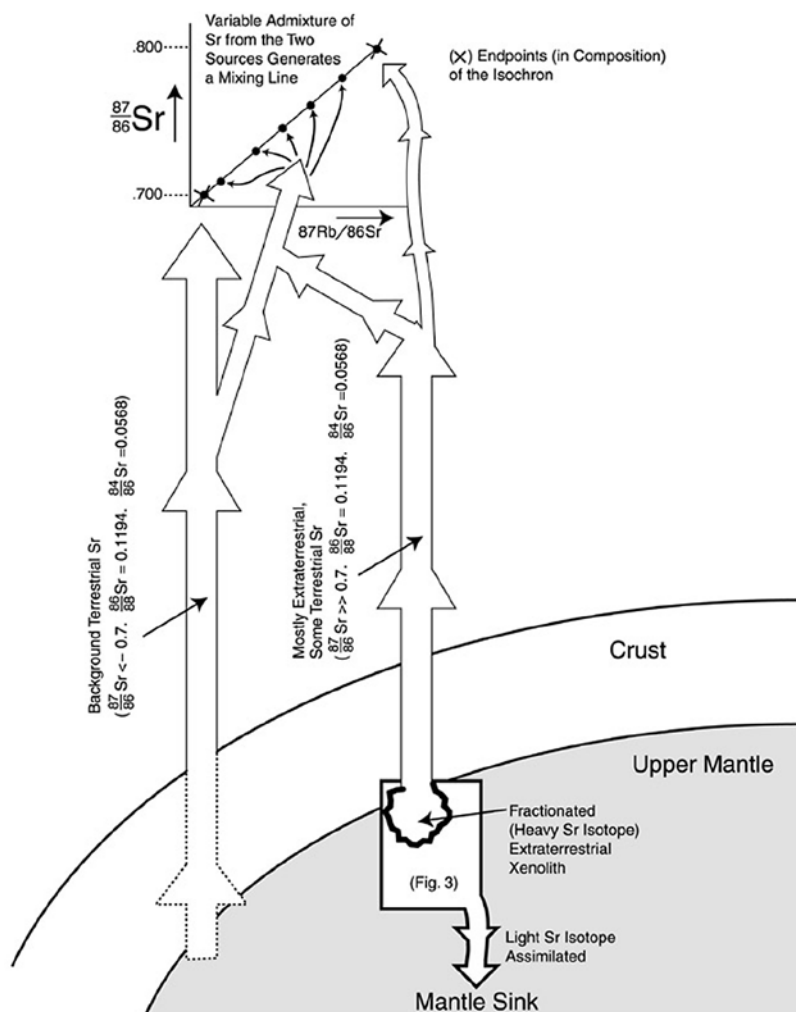


Figure 4. Mixing lines are easily generated by the mixing of material from a fractionated, heavy-Sr isotope xenolith with material from the light-Sr isotope mantle, without prior radiogenic production of ^{87}Sr .

presently-seen relative abundances of ^{88}Sr and ^{84}Sr , in remnants of the extraterrestrial xenoliths, came about as a result of the elevation of the former and diminishing of the latter during isotopic fractionation in the direction of the heavy isotopes. Examples of this can be seen in Table 5. We end up with elevated $^{87}\text{Sr}/^{86}\text{Sr}$ ratios along with near-normal $^{86}\text{Sr}/^{88}\text{Sr}$ and $^{84}\text{Sr}/^{86}\text{Sr}$ ratios.

After this isotope fractionation, the material with elevated $^{87}\text{Sr}/^{86}\text{Sr}$ ratios became incorporated into igneous bodies. As a result of two-component mixing,⁵⁶ pseudoisochrons were formed which have no time significance (Figure 4). For the vast majority of ‘isochrons’, only a modest elevation of $^{87}\text{Sr}/^{86}\text{Sr}$ ratios occurs above 0.7, and thus a relatively small number of fractionations can account for this (Simulation 1, Table 5). However, there are a few mineral ‘isochrons’, which contain minerals with $^{87}\text{Sr}/^{86}\text{Sr}$ ratios in the hundreds or even thousands. But the amount of material with such extremely high $^{87}\text{Sr}/^{86}\text{Sr}$ ratios is, volumetrically speaking, extremely uncommon in the Earth’s

crust. It can be understood as the outcome of rarely-occurring extremely long series of fractionations (Simulation 3, Table 5). As to the belief that such high $^{87}\text{Sr}/^{86}\text{Sr}$ ratios could only have resulted from *in situ* ^{87}Rb decay, we must remember that most mineral ‘isochrons’ are ‘too young’^{22,39} relative to whole-rock ‘isochrons’, which is another way of saying that the $^{87}\text{Sr}/^{86}\text{Sr}$ ratios found in most minerals are *not* commensurate with the concentration of ^{87}Rb in these same minerals.

Significance of the new theory

The earlier belief that a collinear relationship between $^{87}\text{Rb}/^{86}\text{Sr}$ and $^{87}\text{Sr}/^{86}\text{Sr}$ could only be caused by radioactive decay of ^{87}Rb has given way to the recognition that such a relationship can *also* be a mixing line.^{2,51} The theory developed in this paper takes this a step further, and provides a non-radiogenic explanation for the variable and elevated $^{87}\text{Sr}/^{86}\text{Sr}$ ratios observed in rock.

As a result, the need to postulate acceleration of radioactive-decay rates by many orders of magnitudes is eliminated for the Rb-Sr system. The model developed in this paper can also help solve other geologic problems. For instance, the conventional belief that elevations in the $^{87}\text{Sr}/^{86}\text{Sr}$ ratio above about 0.7 are assumed to be the result of either *in situ* or once-present ^{87}Rb decay causes unnecessary interpretative puzzles in isotopic geochemistry,⁵⁷ and these can now be avoided.

Acknowledgements

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References

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2. Woodmorappe, J., *The Mythology of Modern Dating Methods*, Institute for Creation Research, El Cajon, 1999.
3. In fact, in some terrestrial minerals, ^{87}Sr is the most common isotopic species of strontium. For purposes of studying strontium isotope geology, ^{86}Sr has been chosen as the non-radiogenic stable reference isotope with which to compare ^{87}Sr . Hence the oft-quoted $^{87}\text{Sr}/^{86}\text{Sr}$ ratio.
4. For example, Humphreys, D.R., Accelerated nuclear decay: a viable hypothesis? in: Vardiman, L. *et al.* (Eds), *Radioisotopes and the Age of the Earth*, ICR and CRS, California, Missouri, pp. 333–379, 2000.
5. Slight differences in mass between different isotopes cause their fractionation. Although there are other mechanisms by which isotopes can be separated in nature, these are not explored in this work.
6. Hoefs, J., *Stable Isotope Geochemistry*, 4th Edition, Springer-Verlag,

Table 1. Calculated variation of the ratios and abundances of common-earth Sr isotopes with repeated fractionation. (A) shows the variation from isotopically normal to heavy while (B) shows the variation from isotopically normal to light. $\beta = 0.1$.

A. Isotopically Normal to Heavy Strontium

Number of Fractionations	⁸⁷ Sr/ ⁸⁶ Sr	⁸⁷ Sr/ ⁸⁶ Sr	⁸⁷ Sr/ ⁸⁶ Sr	⁸⁷ Sr/ ⁸⁶ Sr	⁸⁷ Sr/ ⁸⁶ Sr	⁸⁷ Sr/ ⁸⁶ Sr	⁸⁷ Sr Abundance	⁸⁶ Sr Abundance	⁸⁷ Sr Abundance	⁸⁶ Sr Abundance
0	11.797143000	8.375253500	0.119399971	14.744429000	0.054795130	0.7099999135	0.7099999135	0.0499999999	0.98600001	0.005400000
1	11.820737284	8.412942141	0.118844481	14.8702990034	0.054375474	0.711710440	0.711710440	0.049904184	0.98219978	0.005354842
2	11.844378741	8.450600380	0.118331987	14.9952095152	0.054354448	0.713484145	0.713484145	0.049808343	0.97841200	0.005314004
3	11.868047518	8.488289982	0.117801878	15.1211492752	0.054138707	0.715244320	0.715244320	0.049712474	0.97443444	0.005274484
4	11.891803453	8.527028713	0.117274145	15.2481870971	0.053921390	0.717050917	0.717050917	0.049614584	0.97087547	0.005232980
5	11.915587240	8.565400342	0.116748775	15.3762718487	0.053705313	0.718839944	0.718839944	0.049520472	0.96712307	0.005187389
10	12.035220711	8.759844144	0.114157021	16.0330143038	0.054434408	0.727832391	0.727832391	0.049040791	0.94855485	0.005182543
20	12.278103059	9.142137027	0.109144842	17.4318349798	0.052559734	0.744217435	0.744217435	0.048079871	0.91233244	0.004795194
50	13.034557375	10.483227410	0.095390449	224.041425195	0.044791428	0.804140795	0.804140795	0.045194234	0.81075448	0.003793445
100	14.404185310	13.121759575	0.074209293	340.385118458	0.038349745	0.910842051	0.910842051	0.040434130	0.044351932	0.002557850
150	15.919707114	14.424383244	0.040885079	517.145100903	0.031739723	1.031701472	1.031701472	0.035804308	0.034089588	0.001717870
200	17.592240358	20.558251450	0.048442249	785.493498473	0.024145473	1.148597702	1.148597702	0.031340125	0.043950219	0.001149987
300	21.482919403	32.209225823	0.031047005	1813.5832103	0.017759974	1.499294422	1.499294422	0.043174439	0.028794434	0.000511427
500	32.035950144	79.042128477	0.012448281	9442.8328384	0.008182049	2.447918954	2.447918954	0.029900314	0.012115598	0.000099131
1000	84.995817794	74.4343999133	0.001339845	4331.751397	0.001178732	8.579079044	8.579079044	0.011349125	0.001322884	0.000001559
3000	4730.8744108	5924812.8502	0.000000149	1.14734E+13	0.000000308	1252.7943219	1252.7943219	0.000211333	0.000000149	8.54444E-14

B. Isotopically Normal to Light Strontium

Number of Fractionations	⁸⁷ Sr/ ⁸⁶ Sr	⁸⁷ Sr/ ⁸⁶ Sr	⁸⁷ Sr/ ⁸⁶ Sr	⁸⁷ Sr/ ⁸⁶ Sr	⁸⁷ Sr/ ⁸⁶ Sr	⁸⁷ Sr/ ⁸⁶ Sr	⁸⁷ Sr Abundance	⁸⁶ Sr Abundance	⁸⁷ Sr Abundance	⁸⁶ Sr Abundance
0	11.797143000	8.375253500	0.119399971	14.744429	0.054795130	0.7099999135	0.7099999135	0.0499999999	0.98600001	0.005400000
-1	11.773595808	8.337733498	0.119934448	14.42359084	0.057015439	0.708172238	0.708172238	0.070095787	0.98981249	0.005443480
-2	11.750095417	8.300381979	0.120474383	14.50177592	0.057237003	0.704409739	0.704409739	0.070191544	0.99943784	0.005487285
-3	11.724442332	8.243197590	0.121018527	143.8097572	0.057459228	0.704451424	0.704451424	0.070287274	0.99974754	0.005731417
-4	11.703235841	8.224179781	0.121543110	142.411818	0.057482315	0.702897889	0.702897889	0.070382974	0.100132574	0.005775879
-5	11.679874109	8.189527804	0.122110144	141.4238574	0.057904248	0.701148514	0.701148514	0.070478444	0.100518832	0.005820472
-10	11.543774883	8.007529553	0.124882441	135.4308553	0.059039145	0.692444745	0.692444745	0.070954474	0.102449143	0.004049491
-20	11.335023195	7.435950897	0.130417348	124.747004	0.041371822	0.475424370	0.475424370	0.071909140	0.104445124	0.004333959
-50	10.475542494	4.691152172	0.149451092	97.04105443	0.048937558	0.424772798	0.424772798	0.074732490	0.119234099	0.008219708
-100	9.440413130	5.345691015	0.187044555	43.88542753	0.083475957	0.553349041	0.553349041	0.079245531	0.143210750	0.011983297
-150	8.742157250	4.270774048	0.234149482	42.04934422	0.101545330	0.488324545	0.488324545	0.083359942	0.170423193	0.017329401
-200	7.911021003	3.412005470	0.295082749	27.47702845	0.123279334	0.431297744	0.431297744	0.084839943	0.201345483	0.024821742
-300	4.478289987	2.177788177	0.459184481	11.99045774	0.181424774	0.334147134	0.334147134	0.090944597	0.270593347	0.049147001
-500	4.344242488	0.887212024	1.127124294	2.250444785	0.394238321	0.204224144	0.204224144	0.082141184	0.402304953	0.158404103
-1000	1.599741744	0.093984439	10.440034434	0.03434392	2.754572849	0.058749147	0.058749147	0.015105301	0.257115234	0.703414580
-3000	0.029417959	0.000011835	84495.887432	1.84284E-09	4353.2924018	0.000402312	0.000402312	0.000000043	0.000157574	0.999842541

Table 2. Calculated variation of the ratios and abundances of common-earth Sr isotopes with repeated fractionation from isotopically normal to heavy, $\beta = 0.05$. Note that the average obtained by mixing equal volumes of materials fractionated from a range of zero to 200 times, is almost identical to the isotopic abundance after 100 continuous episodes of

Number of Fracturines	$^{87}\text{Sr}/^{86}\text{Sr}$	$^{87}\text{Sr}/^{86}\text{Sr}$	$^{87}\text{Sr}/^{86}\text{Sr}$	$^{87}\text{Sr}/^{86}\text{Sr}$	$^{87}\text{Sr}/^{86}\text{Sr}$	$^{87}\text{Sr}/^{86}\text{Sr}$	$^{87}\text{Sr}/^{86}\text{Sr}$	$^{87}\text{Sr}/^{86}\text{Sr}$	$^{87}\text{Sr}/^{86}\text{Sr}$	^{87}Sr Abundance	^{86}Sr Abundance	^{87}Sr Abundance
0	11.797143	8.3752535	0.119599371	147.46429	0.05679513	0.70999135	0.8258	0.069999999	0.098000001	0.0056		
10	11.95947562	8.544274184	0.117037443	154.2361852	0.055397359	0.71563282	0.828310474	0.069975783	0.09604334	0.00570403		
20	12.08332549	8.716703868	0.114722237	161.3190612	0.054033949	0.721371083	0.830786669	0.068753665	0.095309706	0.00514996		
30	12.22931352	8.89261739	0.112452831	168.7271989	0.052704113	0.727153905	0.833228973	0.06813375	0.093698937	0.00493832		
40	12.37686008	9.072078976	0.110228317	176.4753551	0.051407006	0.732987116	0.835657771	0.067516136	0.092110945	0.004755148		
50	12.52618699	9.25516227	0.108047808	184.5796926	0.050141823	0.738863089	0.838013446	0.066900921	0.090545516	0.004540117		
60	12.67731554	9.44194036	0.105910434	193.0360114	0.048907777	0.744790199	0.84035658	0.066288196	0.08902509	0.004352915		
70	12.83026745	9.632487813	0.10381594	201.921582	0.047704102	0.750763823	0.842666954	0.065678051	0.087481737	0.004173239		
80	12.98506472	9.826880896	0.101761891	211.1942797	0.046530051	0.756783343	0.844945546	0.065070569	0.085963088	0.004000797		
90	13.14172963	10.02519662	0.099748667	220.8928007	0.045384895	0.762852143	0.847192532	0.064463832	0.084506326	0.003835311		
100	13.3002847	10.22751474	0.097775464	231.0366999	0.044267923	0.76896961	0.849408285	0.063863918	0.083051289	0.003676508		
110	13.46075274	10.43391594	0.095941294	241.64643	0.04317944	0.775136134	0.851693178	0.0632649	0.081617793	0.003524129		
120	13.62315683	10.64448232	0.093945386	252.743383	0.042115771	0.781352109	0.853747581	0.062668851	0.080205646	0.003377923		
130	13.78730352	10.85929822	0.092086862	264.3499333	0.041079255	0.787617951	0.855671839	0.062079837	0.078814637	0.003237647		
140	13.95386686	11.07844952	0.09036534	276.494827	0.040068249	0.793933999	0.857366379	0.061493923	0.077444627	0.003103071		
150	14.12220257	11.3020231	0.088479734	289.1865078	0.039083125	0.800300718	0.860031503	0.06089917	0.076095339	0.002973968		
160	14.29260307	11.53100882	0.08672945	302.4666092	0.03812027	0.806718493	0.862067589	0.060315638	0.074766648	0.002850125		
170	14.46504546	11.76279752	0.08501379	316.3565629	0.037182088	0.813187753	0.864074996	0.059755881	0.07345829	0.002731333		
180	14.63936635	12.0018211	0.083332069	330.8843749	0.036266995	0.819708851	0.866054078	0.059158452	0.072170078	0.002617392		
190	14.81619283	12.24233734	0.081683615	346.0793371	0.035374434	0.826282263	0.868005187	0.058584901	0.070901801	0.00250811		
200	14.99495051	12.48441991	0.08006777	361.9720864	0.034503821	0.832908389	0.869928673	0.058014775	0.069653249	0.002403303		
Average	13.3002847	10.17238009	0.09830541	225.6606145	0.045481321	0.7659660717	0.848942288	0.063916697	0.08344579	0.003795225		
Compare with 100 Fracturines												
100 Fracturines	13.3002847	10.22751474	0.097775464	231.0366999	0.044267923	0.76896961	0.849408285	0.063863918	0.083051289	0.003676508		

Table 3. Mixtures of assorted combinations of fractionated and unfractionated common-earth Sr. $\beta = 0.5$ (first), and 0.1 (rest). The simulations are intended to approach a final $^{86}\text{Sr}/^{88}\text{Sr}$ ratio of 0.1194. However, it can be seen that mixtures of previously-fractionated strontium cannot be combined in any manner that elevates the $^{87}\text{Sr}/^{86}\text{Sr}$ ratios and simultaneously recovers the near-normal bulk-Earth $^{86}\text{Sr}/^{88}\text{Sr}$ and $^{84}\text{Sr}/^{88}\text{Sr}$ ratios.

Mixing Simulation 1:

Number of Fractions	$^{87}\text{Sr}/^{86}\text{Sr}$	$^{86}\text{Sr}/^{88}\text{Sr}$	$^{84}\text{Sr}/^{88}\text{Sr}$	$^{87}\text{Sr}/^{86}\text{Sr}$	$^{86}\text{Sr}/^{88}\text{Sr}$	$^{84}\text{Sr}/^{88}\text{Sr}$	^{86}Sr Abundance	^{87}Sr Abundance	^{88}Sr Abundance	^{84}Sr Abundance
2000	129 8550118	455 4503325	0 002195629	1171032 727	0 00038893	3 50737585	0 990199619	0 007625425	0 002174111	2 45578E-07
-44E	6 89330515	3 42183249	0 292241094	19 72937186	0 173438479	0 496384956	0 672047063	0 097489291	0 196399769	0 034063277
Average	15 81355368	8 370923111	0 119461138	48 79757776	0 171543824	0 529351168	0 831123341	0 052557658	0 09928694	0 017032061

Mixing Simulation 2:

Number of Fractions	$^{87}\text{Sr}/^{86}\text{Sr}$	$^{86}\text{Sr}/^{88}\text{Sr}$	$^{84}\text{Sr}/^{88}\text{Sr}$	$^{87}\text{Sr}/^{86}\text{Sr}$	$^{86}\text{Sr}/^{88}\text{Sr}$	$^{84}\text{Sr}/^{88}\text{Sr}$	^{86}Sr Abundance	^{87}Sr Abundance	^{88}Sr Abundance	^{84}Sr Abundance
-1000	1 599761764	0 093984639	10 64003663	0 03434392	2 736572869	0 058749147	0 024164883	0 015105301	0 257119236	0 70361458
0	11 797143	8 3752535	0 119399371	147 46429	0 05679513	0 709939135	0 8258	0 069999999	0 098600001	0 0056
0	11 797143	8 3752535	0 119399371	147 46429	0 05679513	0 709939135	0 8258	0 069999999	0 098600001	0 0056
10	12 03522071	8 759864166	0 114157021	100 330163	0 054636408	0 727852391	0 830921161	0 069040791	0 094855485	0 005182563
10	12 03522071	8 759864166	0 114157021	100 330163	0 054636408	0 727852391	0 830921161	0 069040791	0 094855485	0 005182563
10	12 03522071	8 759864166	0 114157021	100 330163	0 054636408	0 727852391	0 830921161	0 069040791	0 094855485	0 005182563
10	12 03522071	8 759864166	0 114157021	100 330163	0 054636408	0 727852391	0 830921161	0 069040791	0 094855485	0 005182563
10	12 03522071	8 759864166	0 114157021	100 330163	0 054636408	0 727852391	0 830921161	0 069040791	0 094855485	0 005182563
1000	86 9982178	746 3439991	0 001339265	633175 1597	0 001178732	8 579079064	0 987326432	0 011349125	0 001322884	1 55933E-06
1000	86 9982178	746 3439991	0 001339265	633175 1597	0 001178732	8 579079064	0 987326432	0 011349125	0 001322884	1 55933E-06
Average	14 92334902	8 38137487	0 11932699	1053692728	0 7954252	0 561625777	0 780502355	0 052300751	0 093123843	0 074073051

Mixing Simulation 3:

Number of Fractions	$^{87}\text{Sr}/^{86}\text{Sr}$	$^{86}\text{Sr}/^{88}\text{Sr}$	$^{84}\text{Sr}/^{88}\text{Sr}$	$^{87}\text{Sr}/^{86}\text{Sr}$	$^{86}\text{Sr}/^{88}\text{Sr}$	$^{84}\text{Sr}/^{88}\text{Sr}$	^{86}Sr Abundance	^{87}Sr Abundance	^{88}Sr Abundance	^{84}Sr Abundance
0	11 797143	8 3752535	0 119399371	147 46429	0 05679513	0 709939135	0 8258	0 069999999	0 098600001	0 0056
0	11 797143	8 3752535	0 119399371	147 46429	0 05679513	0 709939135	0 8258	0 069999999	0 098600001	0 0056
0	11 797143	8 3752535	0 119399371	147 46429	0 05679513	0 709939135	0 8258	0 069999999	0 098600001	0 0056
0	11 797143	8 3752535	0 119399371	147 46429	0 05679513	0 709939135	0 8258	0 069999999	0 098600001	0 0056
0	11 797143	8 3752535	0 119399371	147 46429	0 05679513	0 709939135	0 8258	0 069999999	0 098600001	0 0056
0	11 797143	8 3752535	0 119399371	147 46429	0 05679513	0 709939135	0 8258	0 069999999	0 098600001	0 0056
0	11 797143	8 3752535	0 119399371	147 46429	0 05679513	0 709939135	0 8258	0 069999999	0 098600001	0 0056
0	11 797143	8 3752535	0 119399371	147 46429	0 05679513	0 709939135	0 8258	0 069999999	0 098600001	0 0056
0	11 797143	8 3752535	0 119399371	147 46429	0 05679513	0 709939135	0 8258	0 069999999	0 098600001	0 0056
0	11 797143	8 3752535	0 119399371	147 46429	0 05679513	0 709939135	0 8258	0 069999999	0 098600001	0 0056
100	14 40618531	13 12175937	0 076209293	340 3851185	0 038549745	0 910842051	0 870654088	0 06043613	0 066351932	0 00255785
100	14 40618531	13 12175937	0 076209293	340 3851185	0 038549745	0 910842051	0 870654088	0 06043613	0 066351932	0 00255785
250	19 44049087	25 7329741	0 03886125	1193 702533	0 02158694	1 32865806	0 916474351	0 047142552	0 035615339	0 000767758
Average	12 51748106	9 371320401	0 106707866	173 9920155	0 05386098	0 748663438	0 838684466	0 067001057	0 089494229	0 004820247

Table 5. Fractionation of Sr isotopes whose original ratios were comparable to those modelled to occur as a result of certain stellar nucleosynthetic processes, and which are very different from those encountered in bulk Earth materials and used in Tables 1–4. $\beta = 0.05$.

Reference Values: Bulk Earth Sunium Isotopic Ratios:							Standard Normalized Isotopic Abundances:			
Number of Fissionium	$^{87}\text{Sr}/^{86}\text{Sr}$	$^{87}\text{Sr}/^{84}\text{Sr}$	$^{86}\text{Sr}/^{84}\text{Sr}$	$^{87}\text{Sr}/^{88}\text{Sr}$	$^{86}\text{Sr}/^{88}\text{Sr}$	$^{87}\text{Sr}/^{88}\text{Sr}$	^{87}Sr Abundance	^{86}Sr Abundance	^{84}Sr Abundance	^{88}Sr Abundance
0	11 79714286	2 37525355	0 11939937	147 4642857	0 056795132	0 709939148	0 3258	0 07	0 0926	0 0056
Simulation 1:							Predicted Non-Normalized Abundances:			
							0 5	0 056	0 08	0 004
Altered Isotopic Ratios Before and After Fissionium:							Resulting Normalized Abundances:			
Number of Fissionium	$^{87}\text{Sr}/^{86}\text{Sr}$	$^{87}\text{Sr}/^{84}\text{Sr}$	$^{86}\text{Sr}/^{84}\text{Sr}$	$^{87}\text{Sr}/^{88}\text{Sr}$	$^{86}\text{Sr}/^{88}\text{Sr}$	$^{87}\text{Sr}/^{88}\text{Sr}$	^{87}Sr Abundance	^{86}Sr Abundance	^{84}Sr Abundance	^{88}Sr Abundance
0	2 928571429	6 25	0 16	125	0 05	0 7	0 78125	0 0375	0 125	0 00625
146	10 6371 3703	2 366953607	0 119517813	240 769467	0 034750891	0 78657 947	0 321232862	0 077204314	0 092151956	0 003410268
Simulation 2:							Predicted Non-Normalized Abundances:			
							0 00000001	0 0000007	0 000001	0 0002
Altered Isotopic Ratios Before and After Fissionium:							Resulting Normalized Abundances:			
Number of Fissionium	$^{87}\text{Sr}/^{86}\text{Sr}$	$^{87}\text{Sr}/^{84}\text{Sr}$	$^{86}\text{Sr}/^{84}\text{Sr}$	$^{87}\text{Sr}/^{88}\text{Sr}$	$^{86}\text{Sr}/^{88}\text{Sr}$	$^{87}\text{Sr}/^{88}\text{Sr}$	^{87}Sr Abundance	^{86}Sr Abundance	^{84}Sr Abundance	^{88}Sr Abundance
0	0 014285714	0 01	100	0 00005	200	0 7	4 95761E-05	0 003470329	0 004957612	0 991522483
3368	0 21137039	2 365328542	0 119540174	124 6594663	0 045301704	1 031316305	0 42412547	0 5228777	0 050700032	0 002296792
Simulation 3:							Predicted Non-Normalized Abundances:			
							1E-15	0 0000007	0 000001	100000
Altered Isotopic Ratios Before and After Fissionium:							Resulting Normalized Abundances:			
Number of Fissionium	$^{87}\text{Sr}/^{86}\text{Sr}$	$^{87}\text{Sr}/^{84}\text{Sr}$	$^{86}\text{Sr}/^{84}\text{Sr}$	$^{87}\text{Sr}/^{88}\text{Sr}$	$^{86}\text{Sr}/^{88}\text{Sr}$	$^{87}\text{Sr}/^{88}\text{Sr}$	^{87}Sr Abundance	^{86}Sr Abundance	^{84}Sr Abundance	^{88}Sr Abundance
0	1 42857E-09	0 000000001	1000000000	1E-20	1E+11	0 7	1E-20	7E-12	1E-11	1
11435	0 001290381	2 36364729	0 119565061	198 4138211	0 042152119	6481 531624	0 001282511	0 928550934	0 000154061	6 494E-06

Berlin, Heidelberg, New York, 1997.

- The words 'heavy elements' are used here in an informal sense to refer to elements heavier than about nitrogen. Until fairly recently, spectrometers were not sensitive enough to detect isotopic differences, on a routine basis, for many of the heavy elements. Of course, the published literature is limited to individual studies of individual elements, and these are much more likely to reflect the interest of the researcher than the extent of actual isotopic fractionation in nature.
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- The double-spike technique has to be used in order to detect and measure the extent of the strontium isotope anomalies in both terrestrial and extraterrestrial materials (e.g. Patchett, Ref. 10). In rare cases, to make sure that real natural isotope fractionation effects in strontium are not overlooked (notably within certain meteorites), an automatic correction to the $^{84}\text{Sr}/^{88}\text{Sr}$ of 0.006745 is made (Papanastassiou, D.A. and Wasserburg, G.J., Strontium isotope anomalies in the Allende meteorite, *Geophysical Research Letters* **5**(7):595–598, 1978) in addition to the standard one for $^{86}\text{Sr}/^{88}\text{Sr}$. In the vast majority of investigations, though, the $^{84}\text{Sr}/^{88}\text{Sr}$ ratio is neglected, partly because of the difficulty of measuring it accurately.
- For example, in the Allende meteorite, some inclusions show negative anomalies in ^{142}Nd , ^{146}Nd , ^{148}Nd , and ^{150}Nd . By contrast, the anomaly in ^{145}Nd is positive (McCulloch, M.T. and Wasserburg, G.J., Barium and neodymium isotopic anomalies in the Allende meteorite, *The Astrophysical J.* **220**:L15–L19, 1978). For some elements where a nucleosynthetic origin has been advocated to explain their isotopic anomalies (e.g. samarium: Clayton, D.D., On isotopic anomalies in samarium, *Earth and Planetary Science Letters* **42**:7–12, 1979), a straightforward steady-state nuclear process has been ruled out. Instead, a veritable 'dance' of *ad hoc* stellar processes has been invoked.
- Clayton, R.N., Hinton, R.W. and Davis, A.M., Isotopic variations in the rock-forming elements in meteorites, *Transactions of the Royal Society of London* **A325**: 483–501, 1988.
- Esat, T.M. and Taylor, S.R., Isotope fractionation in the solar system, *International Geology Review* **41**:31–46, 1999.
- For instance, some of the calcium and titanium isotope anomalies in the Allende and Leoville meteorites are consistent with an origin in terms of neutron-rich nucleosynthetic processes in stars, while the absence of certain calcium and titanium isotopes is inconsistent with such an interpretation (Niederer, F.R. *et al.*, Absolute isotopic abundances of Ti in meteorites, *Geochimica et Cosmochimica Acta* **49**:835–851, 1985). Elsewhere, the anomalies in titanium isotopes are instead attributed to neutron-capture s-processes (slow nuclear processes) within stars, with the anomalies scaling in proportion to the known neutron-capture cross sections of the respective titanium isotopes (Ireland, T.R. *et al.*, Isotopically anomalous Ti in presolar SiC from the Murchison meteorite, *The Astrophysical J.* **376**:L53–L56, 1991).

16. Niederer, F.R. and Papanastassiou, D.A., Ca isotopes in refractory inclusions, *Geochimica et Cosmochimica Acta* **48**:1279–1293, 1984.
17. Esat, T.M. *et al.*, Magnesium isotope composition of interplanetary dust particles, *Science* **206**:190–197, 1979.
18. Clayton, D.D., On strontium isotopic anomalies and odd-A, p-process abundances, *The Astrophysical J.* **224**:L93–L95, 1978.
19. Richter, U.O. and Bagemann, F., S-process isotope anomalies: neodymium, samarium, and a bit more of strontium, *Abstracts of Papers to the 23rd Lunar and Planetary Science Conference*, pp. 1147–1148, 1992.
20. The data is compatible with either a 3.2 permil deficit in either ^{84}Sr or ^{88}Sr , or else a 1.6 permil excess of ^{86}Sr (Papanastassiou and Wasserburg, Ref. 11). Exposure to a neutron flux would explain the pattern of Sr isotopes but not the simultaneous existence of large amounts of ^{149}Sm , whose neutron-capture cross-section is much greater than that of the relevant strontium isotopes, and which therefore should have been consumed in a neutron-rich environment.
21. Thirlwall, M.F., Long-term reproducibility of multicollector Sr and Nd isotope ratio analysis, *Chemical Geology* **94**:85–104, 1991. In addition, throughout his lifetime of work on strontium isotopes, one strontium-isotope specialist (personal communication) reports encountering the aforementioned ratio within the range of 0.116–0.123. But, as discussed earlier, it is impossible to know, in the absence of specialized studies, whether any ratio deviating from 0.1194 is caused by natural isotopic fractionation, or if it is an artefact of the spectrometer. It appears that the general belief among scientists has long been that the nonradiogenic isotopes of strontium are identical everywhere on Earth, and that it is unfruitful to search for isotopic anomalies among them.
22. Hidaka, H. and Masuda, A., Nuclide analyses of rare earth elements of the Oklo uranium ore samples: a new method to estimate the neutron fluence, *Earth and Planetary Science Letters* **88**:330–336, 1988.
23. Hidaka, H., Sugiyama, T., Ebihara, M. and Holliger, P., Isotopic evidence for the retention of ^{90}Sr inferred from excess ^{90}Zr in the Oklo natural fission reactors: implication for geochemical behaviour of fissionogenic Rb, Sr, Cs, and Ba, *Earth and Planetary Science Letters* **122**:173–182, 1994. The quoted value is +148 permil isotopically lighter than the usual $^{88}\text{Sr}/^{86}\text{Sr}$ ratio of 8.37.
24. It is interesting to note that some irregularities in uranium isotopes have been known for decades (Cowan, G.A. and Adler, H.H., The variability of the natural abundance of ^{235}U , *Geochimica et Cosmochimica Acta* **40**:1487–1490, 1976), but nothing was made of this fact until decades later, after the existence of natural nuclear reactors had been widely accepted. Cowan and Adler acknowledge that, had the data been published earlier, it would have been attributed to isotope fractionation. Have other terrestrial isotopic anomalies likewise been found, but ignored, because they don't fit current theories? One world-class authority on strontium isotopes suggested that, were an investigator to obtain an anomalous isotope reading, he or she would be prone to regard it as instrument error and ignore it. How many samples of terrestrial materials have been explicitly checked for isotopic anomalies in just the element strontium? I have been unable to find a clear answer in any authoritative source.
25. Morozova, I.M. and Alferovskiy, A.A., Fractionation of lithium and potassium isotopes in geological processes, *Geochemistry International* **11**(1):17–25, 1974.
26. Verbeek, A.A. and Schreiner, G.D.L., Variations in ^{39}K : ^{41}K ratio and movement of potassium in a granite-amphibolite contact region, *Geochimica et Cosmochimica Acta* **31**:2125–2133, 1967.
27. Beard, B.L. and Johnson, C.M., High precision iron isotope measurements of terrestrial and lunar materials, *Geochimica et Cosmochimica Acta* **63**(11/12):1653–1660, 1999. Beard, B.L. *et al.*, Iron isotope biosignatures, *Science* **285**:1889–1892, 1999. Anbar, A.D. *et al.*, Nonbiological fractionation of iron isotopes, *Science* **288**:126–128, 2000.
28. Marechal, C., Telouk, P. and Albarede, F., Precise analysis of copper and zinc isotopic compositions by plasma-source mass spectrometry, *Chemical Geology* **156**:251–273, 1999. The range of isotopic fractionation in copper (9 permil) is greater than that in zinc (just under 1 permil), probably because of the alternate oxidation states in copper. One should not make much of the very small observed range of isotopic fractionation in zinc because only a few samples were analyzed.
29. Allegre, C.J., Isotope geodynamics, *Earth and Planetary Science Letters* **86**:175–203, 1987.
30. Allegre, C.J. and Turcotte, D.I., Implications of a two-component marble-cake mantle, *Nature* **323**:123–127, 1986.
31. Edwards, B.R. and Russell, J.K., Time scales of magmatic processes: new insights from dynamic models for magmatic assimilation, *Geology* **26**(12):1103–1106, 1998.
32. This form of isotopic fractionation is mass-dependent. The lighter isotopes are, by virtue of this fact, slightly more mobile than their heavier counterparts. A series of isotope-fractionation steps in which strontium isotopes are separated is feasible because it is a trace element. More commonly-occurring elements are unlikely to be fractionated, at least in a series of cumulative fractionation events, because there are hardly any locations where they occur in sufficiently low abundance, and to where they could diffuse.
33. Jambon, A., Isotope fractionation: a kinetic model for crystals growing from magmatic melts, *Geochimica et Cosmochimica Acta* **44**:1373–1380, 1980.
34. The exponent, β , is characteristic of self-diffusion in a gas, for which its value is 0.5. In other words, the diffusivity-ratios of gaseous isotopes scale as the square root of the ratios of their respective masses. I have been told that the β value of 0.5 is essentially a holdover from the work of Harold Urey in the 1940s. Urey assumed that mass resonances apply to mass differences of the same isotope. Subsequently, the exponent, β , of 0.5 had been incorrectly assumed to be applicable to isotopic fractionation in magmatic liquids and in aqueous electrolyte solutions. Perhaps this was done for want of any experimental knowledge of the relevant β exponent for isotopic diffusion in liquid-phase situations.
35. Richter, F.M., Liang, Y. and Davis, A.M., Isotope fractionation by diffusion in molten silicates, *Geochimica et Cosmochimica Acta* **63**(18):2853–2861, 1999.
36. Tsuchiyama, A., Kawamura, K., Nakao, T. and Uyeda, C., Isotopic effects on diffusion in MgO melt simulated by the molecular dynamics (MD) method and implications for isotopic mass fractionation in magmatic systems, *Geochimica et Cosmochimica Acta* **58**:3013–3021, 1994.
37. More recent research (unpublished data), suggests that β may be as high as 0.15, at least for the liquid-phase diffusivity of lithium isotopes. The low value of the exponent, β , diminishes the differences between the mass-dependent behaviour of heavy-element isotopes when compared with their light-element counterparts. Were comparative diffusivities to scale on a one-on-one basis with relative atomic masses, the one amu (atomic mass unit) difference between ^{87}Sr and ^{88}Sr would amount to a meagre 1/87, in contrast to the 1/12 difference pertaining to the one amu difference between, say, the carbon isotopes (^{12}C and ^{13}C). Instead, the disparity between the behaviour of heavy elements and lighter ones is reduced. This makes it easier to entertain the possibility of appreciable mass-dependent isotopic fractionation in elements that are heavier than any in which isotopic fractionation has been found thus far.
38. Brady, J.B., Diffusion data for silicate minerals, glasses, and liquids; in: Ahrens, T.J. (Ed.), *Mineral Physics and Crystallography: A Handbook of Physical Constants*, AGU Reference Shelf 2, pp. 269–290, 1995.
39. Fletcher, R.C. and Hofmann, A.W., Simple models of diffusion and combined diffusion-infiltration metasomatism; in: Hoffman, A.W. *et al.*, (Eds), *Geochemical Transport and Kinetics*, Carnegie Institute of Washington Publication 634, pp. 243–259, 1974. In experiments now in progress, there are reportedly even higher values for diffusion coefficients of solute ions in aqueous solutions.
40. Senftle, F.E. and Bracken, J.T., Theoretical effect of diffusion on isotopic abundance ratios in rocks and associated fluids, *Geochimica et Cosmochimica Acta* **7**:61–76, 1955.
41. At the quoted diffusion coefficients applicable to magmas, an element could only diffuse 0.83 cm in one day and 15.9 cm in one year. The high diffusive coefficients in aqueous solutions cause the same element, in the same two time periods, to diffuse at least 2.94 cm and 60 cm

respectively.

42. Consider, for instance, isotopes in an igneous environment. It is possible that, as crystals grow in a magma, there is isotopic fractionation going on at the crystal boundaries, since there is a concentration-gradient between the crystal and the magma. The diffusion rates of strontium out of minerals at high temperatures can be appreciable (Giletti, B.J., Rb and Sr diffusion in alkali feldspars, with implications for cooling histories of rocks, *Geochimica et Cosmochimica Acta* **55**:1331–1343, 1991), and this observation is consistent with such a suggestion. But apparently no such fractionation has ever been measured.
43. That is the main reason for the permil notation: even a percent notation would be inconveniently large in most cases (e.g. 1% equals 10 permil). Perhaps this is also why isotopic fractionation has not been studied in great detail for the heavier elements: it is commonly believed that any isotopic fractionation would be so small that, even if measurable, would be of questionable geologic significance. Moreover, it is supposed, such fractionation would have to be local in nature, as in a thin boundary layer that occurs between rock types or mineral phases of the same rock (Figure 1).
44. Baumgardner, J.R., Distribution of radioactive isotopes in the Earth; in: Vardiman, L. *et al.*, (Eds), *Radioisotopes and the Age of the Earth*, ICR and CRS, California, Missouri, p. 50, 2000. Of course, there were also extensive within-Earth movements during the Flood, but the primary movements within the earth must have been associated with the events which assembled planet Earth during Creation week.
45. As summarized (by Meissner, R. and Wever, Th., The possible role of fluids for the structuring of the continental crust, *Earth-Science Reviews* **32**:19–32, 1992), this includes the unexpected discovery of meteoric water in 20-km boreholes, as well as a variety of indirect evidences for deep crustal water. Much of the lower crust, in fact, may be a saturated environment with excess water in equilibrium with a hydrated mineralogy unfamiliar at the surface, thus explaining the high electrical conductivities which are measured (Gough, D.I., Seismic reflectors, conductivity, water and stress in the continental crust, *Nature* **323**:143–144, 1986).
46. Dreibus, G., Jagoutz, E. and Wanke, H., Water in the Earth's mantle, *Russian Geology and Geophysics* **38**(1):287–293, 1997.
47. Hunter, M.J., The pre-Flood/Flood boundary at the base of the Earth's transition zone, *CEN Tech. J.* **14**(1):70, 2000.
48. McDonough, W.F., Sun, S.-s., Ringwood, A.E., Jagoutz, E. and Hofmann, A.W., Potassium, rubidium, and caesium in the Earth and Moon and the evolution of the mantle of the Earth, *Geochimica et Cosmochimica Acta* **56**:1001–1012, 1992.
49. Russell, W.A., Papanastassiou, D.A. and Tombrello, T.A., Ca isotope fractionation on the Earth and other solar system materials, *Geochimica et Cosmochimica Acta* **42**:1075–1090, 1978.
50. Holden, N.E., Table of the isotopes; in: Lide, D.R., (Ed.), *Handbook of Chemistry and Physics*, 79th Edition, CRC Press, Boca Raton, Boston, London (irregular pagination), 1998.
51. McCulloch, M.T., Primitive ⁸⁷Sr/⁸⁶Sr from an Archean barite and conjecture on the Earth's age and origin, *Earth and Planetary Science Letters* **12**:3–13, 1994.
52. I am not, of course, advocating the uniformitarian idea of the heavy elements forming in stars, but am simply considering the possibility that the processes by which God created the elements and their respective isotopes bears some resemblance to the nucleosynthetic processes believed to occur in stars. This notably involves neutron capture. We can picture God having first created some isotopes by fiat and then having created a shower of energetic neutrons to hit them in order to, in effect, manufacture other isotopes.
53. Fuller, G.M. and Meyer, B.S., Neutrino capture and supernova nucleosynthesis, *Astrophysical J.* **453**:806, 1995. This would occur in a neutrino/antineutrino-rich environment.
54. The, L.-S. *et al.*, A new study of *s*-process nucleosynthesis in massive stars, *Astrophysical J.* **533**:1013, 2000. Some other studies show the same effect.
55. Regardless of temperature (at least in the 5–100 keV range), the neutron-capture cross-section of ⁸⁷Sr is apparently always greater than that of ⁸⁶Sr (Beer, H. *et al.*, On the calculation of Maxwellian-averaged capture cross sections, *Astrophysical J. Supplement Series* **80**:414, 1992). This means that neutron-capture processes will usually convert ⁸⁷Sr to ⁸⁸Sr more rapidly than they convert ⁸⁶Sr to ⁸⁷Sr. However, there is some uncertainty about this conclusion, as there exists an overlap of measurements of neutron-capture cross sections for ⁸⁷Sr and ⁸⁶Sr (Bao, Z.Y. and Kappeler, F., Neutron capture cross sections for *s*-process studies, *Atomic Data and Nuclear Data Tables* **36**:430, 1987).
56. An excellent tutorial showing how two-component mixing produces false isochrons is provided by Helmick, L.S. and Baumann, D.P., A demonstration of the mixing model to account for Rb-Sr isochrons, *Creation Research Society Quarterly* **26**(1)20–23, 1989. Coloured beads are used for each of the relevant isotopes: ⁸⁷Sr, ⁸⁶Sr, and ⁸⁷Rb.
57. For instance, it is unclear whether the measured Rb/Sr ratios are adequate for explaining the ⁸⁷Sr/⁸⁶Sr ratios relevant to the composition of parts of the seafloor (Smith, A.D. and Lewis, C., The planet beyond the plume hypothesis, *Earth-Science Reviews* **48**:135–182, 1999). Moreover, elevated ⁸⁷Sr/⁸⁶Sr signatures are turning up in locations where they are unexpected and where they cannot easily be accounted for by prior crustal contamination (e.g. Prestvik, T., Anomalous Sr and Pb isotope signatures in an off-rift volcano in southeast Iceland, *Geological Society of America Abstracts with Programs* **27**(6)A–108, 1995). A nonradiogenic origin for elevated ⁸⁷Sr/⁸⁶Sr ratios frees researchers to consider other explanations for such isotopic compositions.
58. Richter *et al.*, Ref. 35, Figure 8.

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